

Case study: classical molecular dynamics

- Molecules (e.g. proteins) are modelled by *classical particles*. The motion of the particles is governed by *Newton's 2nd law*:

$$\vec{F} = m \cdot \vec{a}$$

- Since the acceleration $\vec{a}(t) = \dot{\vec{v}}(t)$ is the derivative of the velocity, and the velocity $\vec{v}(t) = \dot{\vec{q}}(t)$ is the derivative of the coordinate $\vec{q}(t)$,
- an MD simulation can be seen as ***solving an ODE numerically***:

$$\ddot{\vec{q}} = \frac{1}{m} \cdot \vec{F}(\vec{q}).$$

MD simulation in practice

- ▶ The ODE is ***discretized***, i.e. transformed to a ***discrete dynamical system***.
- ▶ The standard integration scheme in MD is the ***velocity Verlet algorithm***:

$$\begin{aligned}\mathbf{v}_{n+1/2} &= \mathbf{v}_n + \frac{h}{2m} \mathbf{F}(\mathbf{q}_n) \\ \mathbf{q}_{n+1} &= \mathbf{q}_n + h \cdot \mathbf{v}_{n+1/2} \\ \mathbf{v}_{n+1} &= \mathbf{v}_{n+1/2} + \frac{h}{2m} \mathbf{F}(\mathbf{q}_{n+1}).\end{aligned}$$

- ▶ ***Approximation***: integration is replaced by iteration.
- ▶ ***Crucial parameter***: step size h .

Interpretation of simulation results

The situation

- ▶ The dynamics of typical simulations show sensitive dependence on the initial condition: *Lyapunov instability*.
- ▶ The simulation scheme ***inherits*** this instability.
- ▶ Simulation times are long compared to the Lyapunov time:
- ▶ True orbits and simulated orbits may differ extremely.

Interpretation

- ▶ Not ***true orbits*** are of interest in a simulation but ***statistically defined quantities*** of the system.
- ▶ It is assumed that a simulated ***pseudo orbit*** is good enough,
- ▶ motivated by *shadowing*.

Objections

Quotation

In D. FRENKEL AND B. SMIT, *Understanding Molecular Simulation: From Algorithms to Applications*, 2nd ed., p. 73 we find:

"Hence, our trust in Molecular Dynamics Simulation ... is based largely on belief. To conclude this discussion, let us say that there is clearly still a corpse in the closet. We believe this corpse will not haunt us, and we quickly close the closet."

Exact real arithmetic

- ▶ To support this belief (or to falsify it) it is desirable to simulate true orbits. Here, this is done using the **iRRAM** "machine" (interactive real-RAM, [Mu00]), an implementation of the concept of the **feasible real-RAM** [BH98], which
 1. simulates a RAM register machine,
 2. where each register holds a real number.
 3. The simulation is done by a Turing machine
 4. and works ***iteratively*** with finite,
 5. but ***arbitrary precision***.
- ▶ The iRRAM is a software package, written in C++.
- ▶ The essential data type is `REAL`, a class with an ***arbitrary precision floating point number*** and a fixed precision floating point number representing an ***upper bound on the error*** as core elements.

Example of a discrete dynamical system

Hénon's area-preserving map

Consider the **discrete dynamical system** $f: M \rightarrow M$,

$$x_{n+1} = cx_n - (1 - c)x_n^2 - y_n$$

$$y_{n+1} = x_n - cx_{n+1} + (1 - c)x_{n+1}^2$$

with $M = \mathbb{R}^2$ and control parameter $c \in [-1, 1]$.

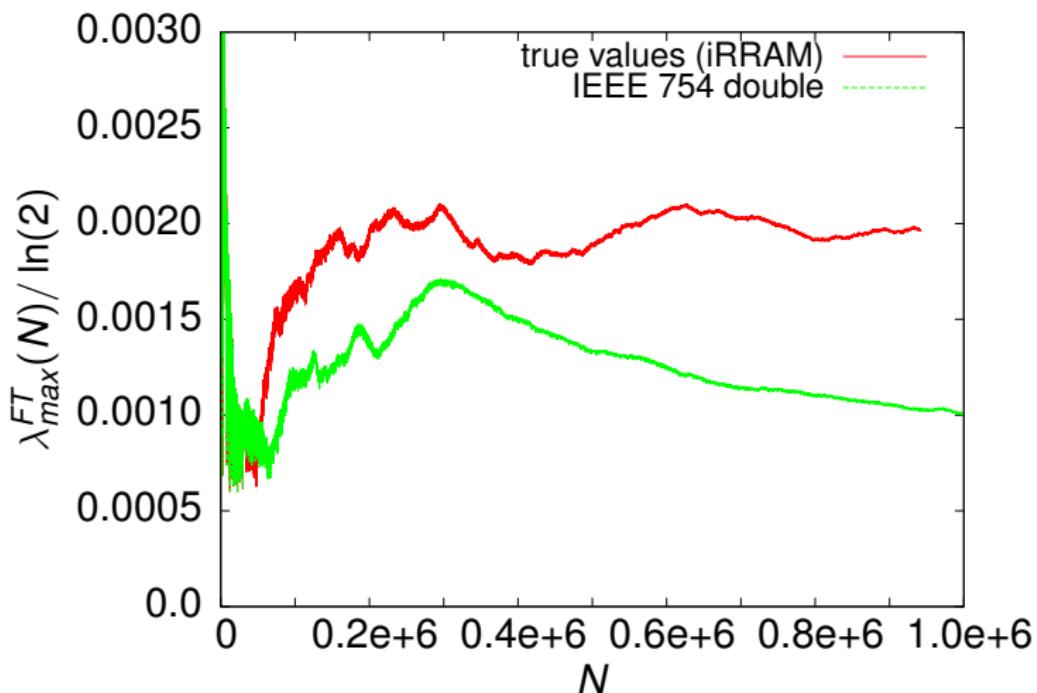
Simulation parameter

$$c = 0.24, \quad x_0 = -0.38484, \quad y_0 = 0.0$$

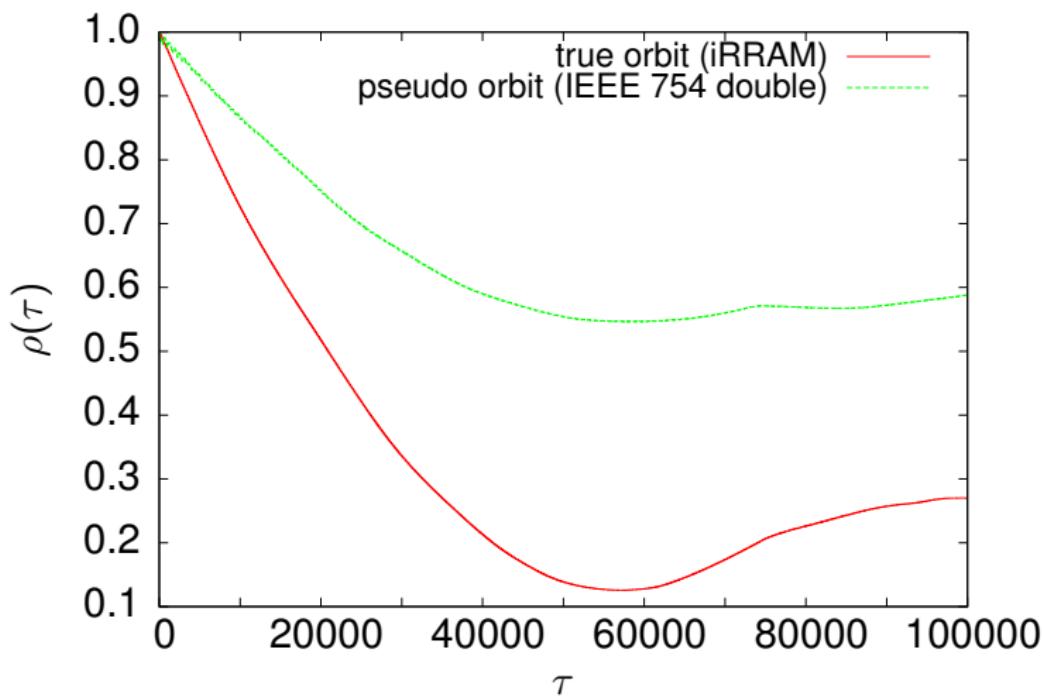
Finite time maximum Lyapunov exponent

$$\lambda_{\max}^{FT}(\mathbf{x}, N) = \frac{1}{N} \cdot \ln \|(Df^N)(\mathbf{x})\|.$$

Finite time maximum Lyapunov exponent



Autocorrelation function



Complexity analysis

Loss of significance

- ▶ The loss of significant bits per iteration, $\sigma(\mathbf{x}, N)$, is 1.4 for the iRRAM in the above example.
- ▶ A theoretical analysis suggests an optimal $\sigma(\mathbf{x}, N)$ **asymptotically** given by the **maximum Lyapunov exponent**:

$$\sigma(\mathbf{x}, N) = \frac{1}{\ln(2)} \cdot \max(0, \lambda_{\max}^{FT}(\mathbf{x}, N))$$

as $N \rightarrow \infty$.

- ▶ The above calculations would expect an optimal loss of significant bits per iteration of 0.002.
- ▶ This would mean in the above example (10^6 iterations): 2000 bits of precision would suffice, instead of 1400000 bits actually used.

For interested persons

The iRRAM C++ package of Norbert Müller can be downloaded at
<https://github.com/norbert-mueller/iRRAM>

Lyapunov instability

Linear stability analysis

- ▶ Consider a small sphere S and apply f N times.
- ▶ The image of S under f^N is **approximated** by **linearization**:

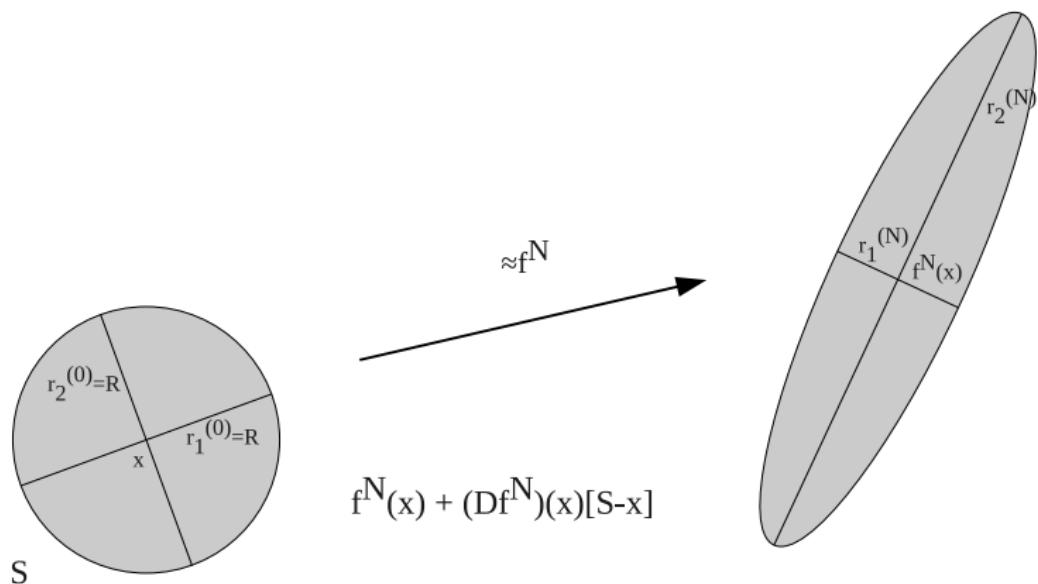
$$f^N(S) \approx f^N(x) + (Df^N)(x) \cdot (S - x).$$

- ▶ The length of the axes of the **ellipsoid** $(Df^N)(x) \cdot (S - x)$ are denoted by $r_1^{(N)}, \dots, r_n^{(N)}$.
- ▶ The value $r_i^{(N)}$ measures the **contraction** or **expansion** of the dynamics.

Lyapunov exponents

The i -th **Lyapunov exponent**: $\lambda_i = \lim_{N \rightarrow \infty} \frac{1}{N} \cdot \ln(r_i^{(N)})$.

Lyapunov instability: a picture



The cocycle / triangularization

Linearized map: *cocycle*

QR-decomposition: Triangularization

$$x^{(k+1)} = f(x^{(k)})$$

$$z^{(k+1)} = (Df)(x^{(k)}) \cdot z^{(k)}$$

$$x^{(0)} \in M, \quad z^{(0)} = \mathbb{1}$$

$$Q^{(k+1)} R^{(k+1)} = (Df)(x^{(k)}) Q^{(k)}$$

$$Q^{(0)} = \mathbb{1}$$

$$\Downarrow$$

$$x^{(k+1)} = f(x^{(k)})$$

$$T^{(k+1)} = R^{(k+1)} \cdot T^{(k)}$$

$$x^{(0)} \in M, \quad T^{(0)} = \mathbb{1}$$

$$z^{(k)} = Q^{(k)} \cdot T^{(k)}$$

Connection to Lyapunov exponents: $T_{ii}^{(N)} \sim e^{\lambda_i \cdot N}$ as $N \rightarrow \infty$

Exact stability analysis (centered forms)

- ▶ Consider *boxes*

$$I = [a_1, b_1] \times \cdots \times [a_n, b_n] \subseteq \mathbb{R}^n.$$

- ▶ Each box $I \in \mathbb{IR}^n$ is uniquely represented by

$$I = x + [-1, 1] \cdot e$$

- ▶ with **center** $x = \text{mid}(I) \in \mathbb{R}^n$ and **extent** $e = \frac{1}{2}\text{wid}(I) \in \mathbb{R}_+^n$.
- ▶ Then (assume $f \in C^2$) the inclusion follows (*centered form*):

$$f(I) \subseteq f(x) + W(I, x) \cdot (I - x)$$

$$W(I, x) = (Df)(x) + \frac{1}{2}(D^2f)(I) \cdot (I - x).$$

- ▶ Finally use Lipschitz bound $|(D^2f)(I) \cdot (I - x)| \leq \|I - x\|_\infty L \in \mathbb{R}^{n \times n}$.

The inclusion cocycle

Dynamics of boxes: *inclusion cocycle* Triangularization

$$x^{(k+1)} = f(x^{(k)})$$

$$z^{(k+1)} = |V(x^{(k)}, e^{(k)})| \cdot z^{(k)}$$

$$x^{(0)} \in M, \quad z^{(0)} = \mathbb{1}$$

$$e^{(k)} = z^{(k)} \cdot e^{(0)}, \quad e^{(0)} \in \mathbb{R}_+^n$$

$$P^{(k+1)} S^{(k+1)} = V(x^{(k)}, e^{(k)}) P^{(k)}$$

$$P^{(0)} = \mathbb{1}$$



$$x^{(k+1)} = f(x^{(k)})$$

$$T^{(k+1)} = |S^{(k+1)}| \cdot T^{(k)}$$

$$x^{(0)} \in M, \quad T^{(0)} = \mathbb{1}$$

$$e^{(k)} = |P^{(k)}| \cdot T^{(k)} \cdot e^{(0)}$$

where $V(x, e) = (Df)(x) + \frac{1}{2} \|e\|_\infty \text{sgn}((Df)(x)) \cdot L$.

From the inclusion cocycle to an algorithm

The computational model

Real numbers $x \in \mathbb{R}$ are represented by a

1. **fixed point approximation** $\hat{x} \in \widehat{\mathbb{R}}$ and an
2. **upper bound** $\bar{e} \in \widehat{\mathbb{R}}_+$ on the **error**.

The iteration algorithm

$$\hat{x}^{(k+1)} = \hat{f}(\hat{x}^{(k)})$$

$$\hat{T}^{(k+1)} = (|\hat{S}^{(k+1)}| + c \cdot U) \cdot \hat{T}^{(k)} + \mathbb{1}$$

$$\hat{x}^{(0)} \in M \cap \widehat{\mathbb{R}}^n(p^s), \quad \hat{T}^{(0)} = \mathbb{1}$$

$$\bar{e}^{(k)} = (|\hat{P}^{(k)}| + d \cdot E) \cdot \hat{T}^{(k)} \cdot \beta^{-p^s}.$$